STUDIA MATHEMATICA

Online First version

Berry–Esseen bound for the Brownian motions on hyperbolic spaces

by

Yuichi Shiozawa (Kyoto)

Abstract. We obtain the uniform convergence rate for the Gaussian fluctuation of the radial part of the Brownian motion on a hyperbolic space. We also show that this result is sharp if the dimension of the hyperbolic space is 2 or general odd. Our approach is based on the repetitive use of the Millson formula and the integration by parts formula.

1. Introduction. We are concerned with the Gaussian fluctuation of the radial part of the Brownian motion on a d-dimensional hyperbolic space. In this note, we obtain the uniform convergence rate for the fluctuation in distribution, the so-called Berry–Esseen bound, together with the sharpness for $d = 2$ and general odd $d \geq 3$.

For $d \geq 2$, let \mathbb{H}^d be the *d*-dimensional hyperbolic space with a pole *o*, and let $d = d_{\mathbb{H}^d}$ be the associated distance function. Let $\Delta = \Delta_{\mathbb{H}^d}$ be the Laplace–Beltrami operator, and $X = (\{X_t\}_{t>0}, \{P_x\}_{x\in \mathbb{H}^d})$ the Brownian motion on \mathbb{H}^d generated by $\Delta/2$. Let $R_t^{(d)} = d(o, X_t)$ $(t \geq 0)$ be the radial process and $P = P_o$. Then by the Itô formula applied to $R_t^{(d)}$ $t^{(a)}_t$, we have

(1.1)
$$
R_t^{(d)} = B_t + \frac{d-1}{2} \int_0^t \coth R_s^{(d)} ds,
$$

where B_t is the Brownian motion on $\mathbb R$ (see, e.g., [\[11,](#page-28-0) Example 3.3.3]). Since $\lim_{t\to\infty} R_t^{(d)} = \infty$, we obtain the law of large numbers:

(1.2)
$$
\lim_{t \to \infty} \frac{R_t^{(d)}}{t} = \frac{d-1}{2}, \quad P\text{-a.s.}
$$

(see also [\[7,](#page-28-1) Section 4.1] and [\[14\]](#page-28-2) for the escape rate). Moreover, we realize

Received 9 September 2023; revised 29 June 2024. Published online *.

²⁰²⁰ Mathematics Subject Classification: Primary 60F05; Secondary 58J65.

Key words and phrases: Berry–Esseen bound, Brownian motion, hyperbolic space, Millson formula.

the limiting behavior of the fluctuation in (1.2) as the central limit theorem:

(1.3)
$$
\lim_{t \to \infty} P\left(\frac{R_t^{(d)} - (d-1)t/2}{\sqrt{t}} \ge x\right) = \Phi(x), \quad x \in \mathbb{R}
$$

(see, e.g., $[2, Corollary 3.1]$ $[2, Corollary 3.1]$ and $[13, Theorem 2.1]$ $[13, Theorem 2.1]$), where

$$
\Phi(x) = \frac{1}{\sqrt{2\pi}} \int_{x}^{\infty} e^{-u^2/2} \, \mathrm{d}u, \quad x \in \mathbb{R}.
$$

In connection with the heat equation in \mathbb{H}^d , Vázquez [\[15,](#page-28-5) (5.5)] noted that the limit in [\(1.3\)](#page-1-0) is uniform in $x \in \mathbb{R}$. Our purpose in this note is to establish the uniform convergence rate in [\(1.3\)](#page-1-0):

THEOREM 1.1. For any $d \geq 2$, there exists a constant $c_1 > 0$ such that

$$
(1.4) \qquad \sup_{x \in \mathbb{R}} \left| P\left(\frac{R_t^{(d)} - (d-1)t/2}{\sqrt{t}} \ge x\right) - \Phi(x) \right| \le \frac{c_1}{\sqrt{t}}, \qquad t \ge 1.
$$

Moreover, if $d = 2$ or if $d \geq 3$ is odd, then there exists a constant $c_2 > 0$ such that $\left(\right)$

(1.5)
$$
P\left(\frac{R_t^{(d)} - (d-1)t/2}{\sqrt{t}} \ge 0\right) - \Phi(0) \ge \frac{c_2}{\sqrt{t}}, \quad t \ge 1.
$$

That is, in these cases, the convergence rate of [\(1.4\)](#page-1-1) is sharp.

Theorem [1.1](#page-1-2) provides the convergence rate $t^{-1/2}$, which is consistent with the standard Berry–Esseen theorem for i.i.d. random variables (see, e.g., [\[6,](#page-28-6) Theorem 3.4.17]). Note that for $d = 1, X$ is the Brownian motion on R started at the origin. Since the radial process $R_t^{(1)}$ $t_t^{(1)}$ is the reflecting Brownian motion, for any $t > 0$ we have

$$
P\left(\frac{R_t^{(1)}}{\sqrt{t}} \ge x\right) = \begin{cases} 1 & (x \le 0), \\ 2\Phi(x) & (x > 0). \end{cases}
$$

In other words, it is unnecessary to study the limiting behavior of the fluctuation as in [\(1.3\)](#page-1-0).

To prove Theorem [1.1,](#page-1-2) we make repetitive use of the Millson formula (see [\(2.4\)](#page-3-0) below) together with the integration by parts formula. By the Millson formula, we reduce the calculation of the distribution under consideration to $d = 2$ or $d = 3$. At present, the validity of [\(1.5\)](#page-1-3) is unavailable for even $d \geq 4$ because of implicit constants in the calculation (see Remark [4.3](#page-22-0) below for details).

Even though we know the matching bound on the transition density function of X (see [\(2.5\)](#page-3-1) below), this bound is insufficient for the estimate of the distribution which would be necessary in the proof of [\(1.4\)](#page-1-1). According to [\(1.1\)](#page-0-1) and [\(1.2\)](#page-0-0), we can also regard $R_t^{(d)}$ $t_t^{(a)}$ as the Brownian motion with

linear drift asymptotically. Hence we might guess that the Girsanov theorem (see, e.g., [\[12,](#page-28-7) Theorem 5.1 (p. 191) and Corollary 5.13 (p. 199)]) is effective for (1.1) ; however, as the so-called Novikov condition fails (see (2.6) below), it is unclear whether the Girsanov theorem is applicable to [\(1.1\)](#page-0-1). Since no information is available about the convergence rate in distribution of [\(1.2\)](#page-0-0), we do not know whether the expression [\(1.1\)](#page-0-1) is applicable for the proof of [\(1.4\)](#page-1-1).

Concerning the law of large numbers [\(1.2\)](#page-0-0), Cammarota–De Gregorio– Macci [\[3,](#page-28-8) Proposition 3.2] and Hirao [\[10,](#page-28-9) Theorem 1.1] proved the large deviation principle for the radial part by using [\(2.5\)](#page-3-1). Moreover, Cammarota–De Gregorio–Macci [\[3,](#page-28-8) Section 3, pp. 1560–1563] proved the moderate deviation principle, together with the exponential decay order of the distribution related to asymptotic normality. Theorem [1.1](#page-1-2) establishes the asymptotic normality in the suitable scaling factor.

Our motivation lies in clarifying how the structures of the volume and spectrum determine the long time behavior of a symmetric Markov process. Anker–Setti [\[1,](#page-28-10) Theorem 2] revealed the concentration behavior in distribution of the Brownian motion on a complete and non-compact Riemannian manifold for which the volume growth rate is exponential and the bottom of the L^2 -spectrum of the Laplacian is strictly positive. In particular, this result characterizes the linear growth rate of the radial part of the Brownian motion in terms of the exponential volume growth rate and the bottom of the L 2 -spectrum. For a Riemannian manifold with a pole, Grigor'yan–Hsu [\[7,](#page-28-1) Theorem 4.1] determined the linear growth rate exactly. Theorem [1.1](#page-1-2) is an attempt to provide the second order asymptotic behavior of the linear growth phenomena for a symmetric Markov process in a quantitative way.

We close this Introduction with some words on the contents and notation. In Section [2,](#page-2-0) we collect basic facts about the Brownian motions on hyperbolic spaces. In Sections [3](#page-5-0) and [4,](#page-13-0) we prove Theorem [1.1](#page-1-2) for odd dimensions and even dimensions, respectively. Two lemmas in elementary calculus are postponed to the Appendix.

For a fixed constant $T > 0$, let $f_1(t)$ and $f_2(t)$ be positive functions defined on [T, ∞). We then write $f_1(t) \sim f_2(t)$ if $f_1(t)/f_2(t) \to 1$ as $t \to \infty$. Let S be a set, and let $g_1(s)$ and $g_2(s)$ be positive functions defined on S. We then write $g_1(s) \leq g_2(s)$ if there exists a constant $c > 0$ such that $g_1(s) \leq cg_2(s)$ for all $s \in S$. We further write $g_1(s) \approx g_2(s)$ if $g_1(s) \lesssim g_2(s)$ and $g_2(s) \lesssim g_1(s)$. We also make the conventions $\sum_{n=1}^0 = 0$ and $\prod_{n=1}^0 = 1$.

2. Preliminaries. In this section, we recall the Millson formula and the estimates of the transition density function of the Brownian motion on \mathbb{H}^d .

4 Y. Shiozawa

Let $d \geq 2$. The d-dimensional hyperbolic space \mathbb{H}^d is a spherically symmetric Riemannian manifold with the Riemannian distance given by

$$
ds^2 = dr^2 + (\sinh r)^2 d\theta^2.
$$

Here $d\theta^2$ is the distance on the $(d-1)$ -dimensional surface $S^{d-1} = \{ \theta \in \mathbb{R}^d \mid \theta \in \mathbb{R}^d \}$ $|\theta| = 1$. We write $\omega_d = 2\pi^{d/2}/\Gamma(d/2)$ for the surface area of S^{d-1} . Let d and dv denote the associated distance function and volume measure, respectively.

Let $X = (\{X_t\}_{t \geq 0}, \{P_x\}_{x \in \mathbb{H}^d})$ be the Brownian motion on \mathbb{H}^d generated by $\Delta/2$, half of the Laplace–Beltrami operator on \mathbb{H}^d . Then there exists a Borel measurable function $p_d(t, x, y) : (0, \infty) \times \mathbb{H}^d \times \mathbb{H}^d \to (0, \infty)$ such that

$$
P_x(X_t \in A) = \int_A p_d(t, x, y) v(\mathrm{d}y), \quad x \in \mathbb{H}^d, t > 0, A \in \mathcal{B}(\mathbb{H}^d).
$$

Namely, $p_d(t, x, y)$ is the transition density function of X. Moreover, there exists a Borel measurable function $q_d(t, r) : (0, \infty) \times (0, \infty) \to (0, \infty)$ such that $p_d(t, x, y) = q_d(t, d(x, y))$ for any $x, y \in \mathbb{H}^d$ and $t > 0$, and

$$
(2.1) \quad P_x(d(x, X_t) \in B) = \omega_d \int_B q_d(t, r) \sinh^{d-1} r \, dr, \quad t > 0, B \in \mathcal{B}([0, \infty))
$$

(see, e.g., [\[3,](#page-28-8) §2.2] for details).

It is known that

(2.2)
$$
q_2(t,r) = \frac{2^{1/2}e^{-t/8}}{(2\pi t)^{3/2}} \int_{r}^{\infty} \frac{se^{-s^2/(2t)}}{(\cosh s - \cosh r)^{1/2}} ds,
$$

(2.3)
$$
q_3(t,r) = \frac{e^{-t/2}}{(2\pi t)^{3/2}} \frac{r}{\sinh r} e^{-r^2/(2t)}
$$

(see, e.g., [\[4,](#page-28-11) Section 5.7] or [\[9,](#page-28-12) Section 2] and references therein). The Millson formula is a recursive relation between the transition density functions:

(2.4)
$$
q_d(t,r) = -\frac{e^{-(d-2)t/2}}{2\pi \sinh r} \frac{\partial q_{d-2}}{\partial r}(t,r), \quad t > 0, r > 0
$$

(see, e.g., [\[5,](#page-28-13) [8\]](#page-28-14), [\[4,](#page-28-11) Section 5.7] and references therein). We also see by [\[4,](#page-28-11) Theorem 5.7.2] that

(2.5)
$$
q_d(t,r) \approx \frac{1}{t^{d/2}} \exp\left(-\frac{(d-1)^2}{8}t - \frac{d-1}{2}r - \frac{r^2}{2t}\right) \times (1+r+t)^{(d-3)/2}(1+r), \quad t > 0, r > 0.
$$

Remark 2.1. As mentioned in the Introduction, the Novikov condition fails for (1.1) , that is,

(2.6)
$$
E\left[\exp\left(\frac{1}{2}\int_{0}^{t}\coth^{2}R_{s}^{(d)}\,\mathrm{d}s\right)\right]=\infty, \quad t>0.
$$

To verify this equality, it is enough to show that

(2.7)
$$
E\left[\int_{0}^{t} \coth^{2} R_{s}^{(d)} ds\right] = \infty, \quad t > 0.
$$

We show it by using (2.5) . Fix $t > 0$. Then

$$
E\left[\int_{0}^{t} \coth^{2} R_{s}^{(d)} ds\right] = \int_{0}^{t} E[\coth^{2} R_{s}^{(d)}] ds
$$

= $\omega_{d} \int_{0}^{t} \left(\int_{0}^{\infty} (\coth^{2} r) q_{d}(s, r) \sinh^{d-1} r dr\right) ds = \omega_{d} I(t).$

By [\(2.5\)](#page-3-1), there exists a constant $c_1 = c_1(t) > 0$ such that for any $s \in [0, t]$ and $r \in [0, \sqrt{s}],$

$$
q_d(s,r) \ge \frac{c_1 e^{-r^2/(2s)}}{s^{d/2}}.
$$

By combining this with

$$
\sinh r \sim r, \quad \coth r \sim \frac{1}{r} \quad (r \to 0),
$$

there exists a constant $c_2 = c_2(t) > 0$ such that for any $s \in [0, t]$ and $r \in [0,\sqrt{s}],$

$$
(\coth^2 r)q_d(s,r)\sinh^{d-1} r \ge \frac{c_2}{r^2} \frac{e^{-r^2/(2s)}}{s^{d/2}} r^{d-1} = \frac{c_2 e^{-r^2/(2s)}}{s^{d/2}} r^{d-3}.
$$

This implies that

$$
I(t) \geq \int_{0}^{t} \left(\int_{0}^{\sqrt{s}} (\coth^2 r) q_d(s, r) \sinh^{d-1} r \, dr \right) ds
$$

\n
$$
\geq c_2 \int_{0}^{t} \left(\int_{0}^{\sqrt{s}} e^{-r^2/(2s)} r^{d-3} \, dr \right) \frac{1}{s^{d/2}} ds
$$

\n
$$
= c_2 \int_{0}^{t} \frac{1}{s} ds \left(\int_{0}^{1} e^{-u^2/2} u^{d-3} \, du \right) = \infty.
$$

At the first equality above, we used the change of variables formula with $\subset \mathbb{R}^n$. $r = \sqrt{s} u$. We thus get (2.7) .

Before the proof of Theorem [1.1](#page-1-2) in the subsequent sections, we make a comment on the range of $x \in \mathbb{R}$ in [\(1.4\)](#page-1-1). Since $R_t^{(d)} \geq 0$ for any $t \geq 0$, we comment on the range or $x \in \mathbb{R}$
have, for any $x \leq -(d-1)\sqrt{t/2}$,

$$
P\left(\frac{R_t^{(d)} - (d-1)t/2}{\sqrt{t}} \ge x\right) = 1,
$$

which implies that

$$
\left| P\left(\frac{R_t^{(d)} - (d-1)t/2}{\sqrt{t}} \ge x\right) - \Phi(x) \right| = 1 - \Phi(x)
$$

$$
\le 1 - \Phi\left(-\frac{d-1}{2}\sqrt{t}\right) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{-(d-1)\sqrt{t}/2} e^{-y^2/2} dy.
$$

We also note that for $u > 0$,

(2.8)
$$
\int_{-\infty}^{-u} e^{-y^2/2} dy = \int_{u}^{\infty} e^{-y^2/2} dy = \int_{u}^{\infty} \frac{1}{y} y e^{-y^2/2} dy
$$

$$
\leq \frac{1}{u} \int_{u}^{\infty} y e^{-y^2/2} dy = \frac{e^{-u^2/2}}{u},
$$

which yields

$$
\int_{-\infty}^{-(d-1)\sqrt{t}/2} e^{-y^2/2} dy \le \frac{2e^{-(d-1)^2t/8}}{(d-1)\sqrt{t}}.
$$

Then there exists a constant $c_1 > 0$ such that for any $t \geq 1$,

$$
\sup_{x \le -(d-1)\sqrt{t}/2} \left| P\left(\frac{R_t^{(d)} - (d-1)t/2}{\sqrt{t}} \ge x\right) - \Phi(x) \right| \le \frac{c_1 e^{-(d-1)^2 t/8}}{\sqrt{t}} \le \frac{c_1}{\sqrt{t}}.
$$

Hence for the proof of [\(1.4\)](#page-1-1) it is enough to show that, for some $c_2 > 0$,

(2.9)
$$
\sup_{x \ge -(d-1)\sqrt{t}/2} \left| P\left(\frac{R_t^{(d)} - (d-1)t/2}{\sqrt{t}} \ge x\right) - \Phi(x) \right| \le \frac{c_2}{\sqrt{t}}, \quad t \ge 1.
$$

3. Proof of Theorem [1.1](#page-1-2) for odd dimensions. In this section, we first prove Theorem [1.1](#page-1-2) for $d = 3$. Using this assertion, we next prove the theorem for general odd dimensions.

3.1. Proof of Theorem [1.1](#page-1-2) for $d = 3$. We first prove [\(2.9\)](#page-5-1) for $d = 3$. We write $R_t = R_t^{(3)}$ We write $R_t = R_t^{(3)}$ for simplicity. Since $R_t \ge 0$ for any $t \ge 0$ and $\omega_3 = 4\pi$, it follows by [\(2.3\)](#page-3-3) that for any $x \ge -\sqrt{t}$,

$$
P\left(\frac{R_t - t}{\sqrt{t}} \ge x\right) = P(R_t \ge t + x\sqrt{t})
$$

=
$$
\frac{\omega_3}{(2\pi t)^{3/2}} \int_{t+x\sqrt{t}}^{\infty} e^{-t/2} e^{-u^2/(2t)} u \sinh u \, du
$$

=
$$
\frac{1}{\sqrt{2\pi}} \frac{2e^{-t/2}}{t\sqrt{t}} \int_{t+x\sqrt{t}}^{\infty} e^{-u^2/(2t)} u \sinh u \, du.
$$

Then by the change of variables $u = t + v$ √ t we obtain

$$
\frac{2e^{-t/2}}{t\sqrt{t}}\int_{t+x\sqrt{t}}^{\infty} e^{-u^2/(2t)}u\sinh u \,du
$$

= $2e^{-t}\int_{x}^{\infty} e^{-v^2/2}e^{-v\sqrt{t}}\sinh(t+v\sqrt{t})\,dv$
+ $\frac{2e^{-t}}{\sqrt{t}}\int_{x}^{\infty} e^{-v^2/2}e^{-v\sqrt{t}}v\sinh(t+v\sqrt{t})\,dv$
= $I_1(t,x) + I_2(t,x)$.

Therefore,

$$
(3.1) \ \ P\left(\frac{R_t - t}{\sqrt{t}} \ge x\right) - \Phi(x) = \frac{1}{\sqrt{2\pi}} \Big\{ \Big(I_1(t, x) - \int_x^{\infty} e^{-v^2/2} dv\Big) + I_2(t, x) \Big\}.
$$

We have

(3.2)
$$
I_1(t,x) - \int_x^{\infty} e^{-v^2/2} dv = \int_x^{\infty} e^{-v^2/2} (1 - e^{-2(t+v\sqrt{t})}) dv - \int_x^{\infty} e^{-v^2/2} dv
$$

$$
= -\int_x^{\infty} e^{-v^2/2} e^{-2(t+v\sqrt{t})} dv.
$$

Then by the change of variables $u = v + 2\sqrt{t}$ we get, for any $t \ge 1$ and $x \ge -\sqrt{t}$,

$$
(3.3) \qquad \int\limits_x^\infty e^{-v^2/2} e^{-2(t+v\sqrt{t})} dv = \int\limits_{x+2\sqrt{t}}^\infty e^{-u^2/2} du \le \int\limits_{\sqrt{t}}^\infty e^{-u^2/2} du \le \frac{e^{-t/2}}{\sqrt{t}}.
$$

In the last inequality above, we have used [\(2.8\)](#page-5-2). Hence for any $t \geq 1$ and m the las
 $x \geq -\sqrt{t},$

(3.4)
$$
0 \ge I_1(t,x) - \int_x^{\infty} e^{-u^2/2} du \ge -\frac{e^{-t/2}}{\sqrt{t}}.
$$

Since $(e^{-v^2/2})' = -ve^{-v^2/2}$, by integration by parts we have

(3.5)
$$
I_2(t,x) = \frac{1}{\sqrt{t}} \int_x^{\infty} v e^{-v^2/2} (1 - e^{-2(t+v\sqrt{t})}) dv
$$

$$
= \frac{e^{-x^2/2}}{\sqrt{t}} (1 - e^{-2(t+x\sqrt{t})}) + 2 \int_x^{\infty} e^{-v^2/2} e^{-2(t+v\sqrt{t})} dv \ge 0.
$$

8 Y. Shiozawa

In particular, we see by [\(3.3\)](#page-6-0) that for any $t \geq 1$ and $x \geq -\sqrt{t}$,

$$
0 \le I_2(t, x) \le \frac{e^{-x^2/2}}{\sqrt{t}} + \frac{2e^{-t/2}}{\sqrt{t}} \le \frac{3}{\sqrt{t}}.
$$

Combining this with (3.1) and (3.4) , we get (2.9) for $d = 3$.

We next prove [\(1.5\)](#page-1-3) for $d = 3$. If we take $x = 0$ in [\(3.1\)](#page-6-1), then by a calculation similar to (3.2) and (3.3) , it follows that

$$
I_1(t,0) - \int_0^\infty e^{-u^2/2} \, \mathrm{d}u = -\int_{2\sqrt{t}}^\infty e^{-u^2/2} \, \mathrm{d}u \ge -\frac{e^{-2t}}{2\sqrt{t}} \ge -\frac{e^{-2}}{2\sqrt{t}}, \quad t \ge 1.
$$

We also have, for any $t \geq 1$,

$$
I_2(t,0) = \frac{1}{\sqrt{t}} \int_0^\infty v e^{-v^2/2} (1 - e^{-2(t+v\sqrt{t})}) dv
$$

$$
\geq \frac{1 - e^{-2}}{\sqrt{t}} \int_0^\infty v e^{-v^2/2} dv = \frac{1 - e^{-2}}{\sqrt{t}}.
$$

Hence by [\(3.1\)](#page-6-1) with $x = 0$, for any $t \ge 1$ we get

$$
P\left(\frac{R_t - t}{\sqrt{t}} \ge 0\right) - \Phi(0) = \frac{1}{\sqrt{2\pi}} \left\{ \left(I_1(t, 0) - \int_0^\infty e^{-u^2/2} du \right) + I_2(t, 0) \right\}
$$

$$
\ge \frac{1}{\sqrt{2\pi}} \left(-\frac{e^{-2}}{2\sqrt{t}} + \frac{1 - e^{-2}}{\sqrt{t}} \right) = \frac{2 - 3e^{-2}}{2\sqrt{2\pi}} \frac{1}{\sqrt{t}},
$$

which implies [\(1.5\)](#page-1-3) for $d = 3$.

3.2. Proof of [\(1.4\)](#page-1-1) for odd dimensions. Let $d \geq 2$. For $t \geq 1$ and 5.2. Proof of (1.
 $x \ge -(d-1)\sqrt{t}/2$, let

$$
T = T(t, x) = x\sqrt{t} + \frac{d-1}{2}t \quad (\ge 0).
$$

Since $R_t^{(d)} \geq 0$, by (2.1) and (2.4) we have

(3.6)
$$
P\left(\frac{R_t^{(d)} - (d-1)t/2}{\sqrt{t}} \ge x\right) = \omega_d \int_T^{\infty} q_d(t, r) \sinh^{d-1} r dr
$$

$$
= -\frac{\omega_d e^{-(d-2)t/2}}{2\pi} \int_T^{\infty} \frac{\partial q_{d-2}}{\partial r} (t, r) \sinh^{d-2} r dr.
$$

Then by integration by parts and [\(2.5\)](#page-3-1),

$$
-\int_{T}^{\infty} \frac{\partial q_{d-2}}{\partial r}(t,r) \sinh^{d-2} r \, dr
$$

= $[-q_{d-2}(t,r) \sinh^{d-2} r]_{r=T}^{r=\infty} + \int_{T}^{\infty} q_{d-2}(t,r) \left(\frac{\partial}{\partial r} \sinh^{d-2} r\right) dr$
= $q_{d-2}(t,T) \sinh^{d-2} T + \int_{T}^{\infty} q_{d-2}(t,r) \left(\frac{\partial}{\partial r} \sinh^{d-2} r\right) dr.$

In the same way, we have

$$
\int_{T}^{\infty} q_{d-2}(t,r) \left(\frac{\partial}{\partial r} \sinh^{d-2} r \right) dr
$$
\n
$$
= -\frac{e^{-(d-4)t/2}}{2\pi} \int_{T}^{\infty} \frac{1}{\sinh r} \frac{\partial q_{d-4}}{\partial r} (t,r) \left(\frac{\partial}{\partial r} \sinh^{d-2} r \right) dr
$$
\n
$$
= \frac{e^{-(d-4)t/2}}{2\pi} q_{d-4}(t,T) \left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right) \sinh^{d-2} r \Big|_{r=T}
$$
\n
$$
+ \frac{e^{-(d-4)t/2}}{2\pi} \int_{T}^{\infty} q_{d-4}(t,r) \frac{\partial}{\partial r} \left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right) \sinh^{d-2} r dr.
$$

Repeating this procedure, for any $n \geq 2$ with $d \geq 2n$ we see from [\(3.6\)](#page-7-0) that

$$
(3.7)
$$
\n
$$
P\left(\frac{R_t^{(d)} - (d-1)t/2}{\sqrt{t}} \ge x\right)
$$
\n
$$
= \omega_d \sum_{m=1}^{n-1} \prod_{j=1}^m \frac{e^{-(d-2j)t/2}}{2\pi} q_{d-2m}(t, T) \left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{m-1} \sinh^{d-2} r \Big|_{r=T}
$$
\n
$$
+ \omega_d \prod_{j=1}^{n-1} \frac{e^{-(d-2j)t/2}}{2\pi} \int_{T}^{\infty} \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{n-1} \sinh^{d-2} r\right) q_{d-2n+2}(t, r) \sinh r \, dr
$$
\n
$$
= \omega_d \sum_{m=1}^{n-1} \frac{e^{-(d-(m+1))mt/2}}{(2\pi)^m} q_{d-2m}(t, T) \left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{m-1} \sinh^{d-2} r \Big|_{r=T}
$$
\n
$$
+ \omega_d \frac{e^{-(d-n)(n-1)t/2}}{(2\pi)^{n-1}} \int_{T}^{\infty} \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{n-1} \sinh^{d-2} r\right) q_{d-2n+2}(t, r) \sinh r \, dr.
$$

10 Y. Shiozawa

We now define $J_1^{(2n+1)}$ $J_1^{(2n+1)}(t, x)$ and $J_2^{(2n+1)}$ $2^{(2n+1)}(t, x)$ by $(3.7)₂$ $(3.7)₂$ and $(3.7)₄$, respectively, with $d = 2n + 1$:

$$
(3.8) \quad J_1^{(2n+1)}(t,x) = \omega_{2n+1} \sum_{m=1}^{n-1} \frac{e^{-(2n-m)mt/2}}{(2\pi)^m} q_{2n+1-2m}(t,T) \left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{m-1} \sinh^{2n-1} r\Big|_{r=T},
$$

(3.9)
$$
J_2^{(2n+1)}(t, x) = \omega_{2n+1} \frac{e^{-(n^2-1)t/2}}{(2\pi)^{n-1}} \int_{T}^{\infty} \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{n-1} \sinh^{2n-1} r \right) q_3(t, r) \sinh r \, dr.
$$

Then

(3.10)
\n
$$
P\left(\frac{R_t^{(2n+1)} - nt/2}{\sqrt{t}} \ge x\right) - \Phi(x) = J_1^{(2n+1)}(t, x) + (J_2^{(2n+1)}(t, x) - \Phi(x)).
$$

Let us give an upper bound of $J_1^{(2n+1)}$ $1^{(2n+1)}(t, x).$

LEMMA 3.1. There exists a constant $c > 0$ such that

(3.11)
$$
\sup_{x \ge -n\sqrt{t}} J_1^{(2n+1)}(t, x) \le \frac{c}{\sqrt{t}}, \quad t \ge 1.
$$

Proof. We give an upper bound of $e^{-(d-m-1)mt/2}q_{d-2m}(t, T)$; recall that $d = 2n + 1$. For $m = 1, ..., n - 1$, we see by [\(2.5\)](#page-3-1) that

$$
(3.12) \quad e^{-(d-m-1)mt/2} q_{d-2m}(t,T)
$$
\n
$$
\asymp \frac{1}{t^{(d-2m)/2}} e^{-(d-m-1)mt/2}
$$
\n
$$
\asymp \exp\left(-\frac{(d-2m-1)^2}{8}t - \frac{(d-2m-1)T}{2} - \frac{T^2}{2t}\right)
$$
\n
$$
\asymp (1+T+t)^{(d-2m-3)/2}(1+T).
$$

Then

$$
(3.13) \qquad e^{-(d-m-1)mt/2} \exp\left(-\frac{(d-2m-1)^2}{8}t - \frac{(d-2m-1)T}{2} - \frac{T^2}{2t}\right) \\
= e^{-(d-m-1)T}e^{-x^2/2}.
$$

We can also see that there exists a constant $c_1 > 0$ such that for any $t \geq 1$

and
$$
x \ge -(d-1)\sqrt{t}/2
$$
,
\n $(1+T+t)^{(d-2m-3)/2}(1+T)$
\n $\le (1+T+t)^{(d-2m-1)/2} = (1+x\sqrt{t} + (d-1)t/2 + t)^{(d-2m-1)/2}$
\n $\le c_1(1+(|x|\sqrt{t})^{(d-2m-1)/2} + t^{(d-2m-1)/2}).$

Hence

$$
(3.14) \quad e^{-(d-m-1)mt/2} q_{d-2m}(t,T)
$$
\n
$$
\leq \frac{c_2}{t^{(d-2m)/2}} e^{-(d-m-1)T} e^{-x^2/2} \left(1 + (|x|\sqrt{t})^{(d-2m-1)/2} + t^{(d-2m-1)/2}\right)
$$
\n
$$
\leq c_3 e^{-(d-m-1)T} \left(\frac{1}{t^{(d-2m)/2}} + \frac{e^{-x^2/2} |x|^{(d-2m-1)/2}}{t^{(d-2m)/4+1/4}} + \frac{1}{\sqrt{t}}\right) \leq \frac{c_4 e^{-(d-m-1)T}}{\sqrt{t}}.
$$

At the last inequality, we used the fact that $\sup_{x \in \mathbb{R}} e^{-x^2/2} |x|^{(d-2m-1)/2} < \infty$.

By [\(A.6\)](#page-26-0), for any m ($1 \le m \le n-1$), there exists a constant $c_5 > 0$ such that

$$
\left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{m-1} \sinh^{d-2} r \bigg|_{r=T} \le c_5 e^{(d-m-1)T}
$$

.

By combining this with [\(3.14\)](#page-10-0), there exists a constant $c_6 > 0$ such that for by combining this with (3.14), then
any $t \ge 1$ and $x \ge -(d-1)\sqrt{t}/2$,

$$
(3.15) \quad \frac{e^{-(d-m-1)mt/2}}{(2\pi)^m} q_{d-2m}(t,T) \left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{m-1} \sinh^{d-2} r \Big|_{r=T}
$$

$$
\leq \frac{c_6 e^{-(d-m-1)T}}{\sqrt{t}} \cdot e^{(d-m-1)T} = \frac{c_6}{\sqrt{t}}.
$$

Therefore, by the definition of $J_1^{(2n+1)}$ $1^{(2n+1)}(t, x)$ in $(3.7)_2$ $(3.7)_2$, we arrive at (3.11) .

We turn to the calculation of $J_2^{(2n+1)}$ $2^{(2n+1)}(t, x).$

LEMMA 3.2. Let $n \geq 1$.

(i) For any $t \geq 1$ and $x \geq -n$ √ t,

$$
J_2^{(2n+1)}(t,x) = P\bigg(\frac{R_{n^2t}^{(3)} - n^2t}{\sqrt{n^2t}} \ge x\bigg).
$$

(ii) There exists a constant $c > 0$ such that

(3.16)
$$
\sup_{x \ge -n\sqrt{t}} |J_2^{(2n+1)}(t,x) - \Phi(x)| \le \frac{c}{\sqrt{t}}, \quad t \ge 1.
$$

Proof. We first prove (i). Let $d = 2n + 1$. For $t \ge 1$ and $x \ge -n$ √ st prove (i). Let $d = 2n + 1$. For $t \ge 1$ and $x \ge -n\sqrt{t}$, we then have $T = x\sqrt{t} + nt$. By [\(2.3\)](#page-3-3) and [\(A.9\)](#page-26-1), it follows that

12 Y. Shiozawa

(3.17)
$$
\begin{aligned} \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{n-1} \sinh^{2n-1} r \right) q_3(t, r) \sinh r \\ &= \frac{(2n-1)!!}{n} \sinh(nr) \frac{e^{-t/2}}{(2\pi t)^{3/2}} \frac{r}{\sinh r} e^{-r^2/(2t)} \sinh r \\ &= \frac{(2n-1)!!}{n} \frac{e^{-t/2}}{(2\pi t)^{3/2}} r e^{-r^2/(2t)} \sinh(nr) \end{aligned}
$$

and thus

$$
\frac{e^{-(n^2-1)t/2}}{(2\pi)^{n-1}} \int_{T}^{\infty} \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{n-1} \sinh^{2n-1} r \right) q_3(t, r) \sinh r \, dr
$$

=
$$
\frac{e^{-(n^2-1)t/2}}{(2\pi)^{n-1}} \frac{(2n-1)!!}{n} \frac{e^{-t/2}}{(2\pi t)^{3/2}} \int_{T}^{\infty} r e^{-r^2/(2t)} \sinh(nr) \, dr
$$

=
$$
\frac{1}{(2\pi)^{n-1} (2\pi t)^{3/2}} \frac{(2n-1)!!}{n} \int_{T}^{\infty} r e^{-n^2 t/2} e^{-r^2/(2t)} \sinh(nr) \, dr.
$$

We also note that

$$
\omega_{2n+1} = \frac{(2\pi)^{n-1} 4\pi}{(2n-1)!!} = \frac{(2\pi)^{n-1} \omega_3}{(2n-1)!!}.
$$

Then by (3.9) and (2.3) ,

$$
J_2^{(2n+1)}(t,x) = \frac{\omega_{2n+1}}{(2\pi)^{n-1}(2\pi t)^{3/2}} \frac{(2n-1)!!}{n} \int_T^{\infty} r e^{-n^2 t/2} e^{-r^2/(2t)} \sinh(nr) dr
$$

\n
$$
= \frac{\omega_3}{n(2\pi t)^{3/2}} \int_T^{\infty} r e^{-n^2 t/2} e^{-r^2/(2t)} \sinh(nr) dr
$$

\n
$$
= \frac{\omega_3}{(2\pi n^2 t)^{3/2}} \int_{x\sqrt{n^2 t} + n^2 t}^{\infty} \frac{u}{\sinh u} e^{-n^2 t/2} e^{-u^2/(2n^2 t)} \sinh^2 u du
$$

\n
$$
= \omega_3 \int_{x\sqrt{n^2 t} + n^2 t}^{\infty} q_3(n^2 t, u) \sinh^2 u du
$$

\n
$$
= P\left(\frac{R_{n^2 t}^{(3)} - n^2 t}{\sqrt{n^2 t}} \ge x\right).
$$

At the second equality above, we have used the change of variables $u = nr$. We next prove (ii). By (i),

(3.18)
$$
J_2^{(2n+1)}(t,x) - \Phi(x) = P\left(\frac{R_{n^2t}^{(3)} - n^2t}{\sqrt{n^2t}} \ge x\right) - \Phi(x).
$$

Since [\(1.4\)](#page-1-1) is already proved for $d = 3$ in Section [3.1,](#page-5-3) the proof is complete.

We are now in a position to complete the proof of (1.4) . By (3.10) , (3.11) and [\(3.16\)](#page-10-1), we have

$$
\sup_{x \ge -n\sqrt{t}} \left| P\left(\frac{R_t^{(2n+1)} - nt}{\sqrt{t}} \ge x\right) - \Phi(x) \right| \lesssim \frac{1}{\sqrt{t}}, \quad t \ge 1.
$$

The proof of (1.4) is complete for general odd dimensions.

3.3. Proof of [\(1.5\)](#page-1-3) for odd dimensions. We start from [\(3.10\)](#page-9-2) with $x = 0$. Since $T = T(t, 0) = nt$, we see by [\(3.8\)](#page-9-3) that

$$
J_1^{(2n+1)}(t,0)
$$

= $\omega_{2n+1} \sum_{m=1}^{n-1} \frac{e^{-(2n-m)m t/2}}{(2\pi)^m} q_{2n+1-2m}(t, nt) \left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{m-1} \sinh^{2n-1} r \Big|_{r=nt}.$

By [\(3.12\)](#page-9-4) with $d = 2n + 1$ and $T = nt$, we have

$$
e^{-(2n-m)mt/2}q_{2n+1-2m}(t,nt)
$$

\n
$$
\asymp \frac{1}{t^{(2n+1-2m)/2}} \times e^{-(2n-m)mt/2} \exp\left(-\frac{(n-m)^2}{2}t - (n-m)nt - \frac{n^2t}{2}\right)
$$

\n
$$
\times (1+nt+t)^{n-m-1}(1+nt).
$$

Then by [\(3.13\)](#page-9-5) with $d = 2n + 1$ and $T = nt$,

$$
e^{-(2n-m)mt/2} \exp\left(-\frac{(n-m)^2}{2}t - (n-m)nt - \frac{n^2t}{2}\right) = e^{-(2n-m)nt}.
$$

For any $t \geq 1$, since $(1 + nt + t)^{n-m-1}(1 + nt) \asymp t^{n-m}$, we get

$$
e^{-(2n-m)mt/2}q_{2n+1-2m}(t, nt) \approx \frac{e^{-(2n-m)nt}}{\sqrt{t}}.
$$

On the other hand, it follows by [\(A.7\)](#page-26-2) that for any m ($1 \leq m \leq n-1$), there exists a constant $c_1 > 0$ such that

$$
\left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{m-1} \sinh^{2n-1} r \bigg|_{r=nt} \ge c_1 e^{(2n-m)nt}.
$$

Hence by the argument above,

(3.19)
$$
J_1^{(2n+1)}(t,0) \gtrsim \frac{1}{\sqrt{t}}, \quad t \ge 1.
$$

Since (1.5) is already proved for $d = 3$, we see from (3.18) that

$$
J_2^{(2n+1)}(t,0)-\Phi(0)=P\bigg(\frac{R_{n^2t}^{(3)}-n^2t}{\sqrt{n^2t}}\geq0\bigg)-\Phi(0)\gtrsim\frac{1}{\sqrt{t}},\quad \ \, t\geq1.
$$

Combining this with (3.10) (for $x = 0$) and (3.19) , we have

$$
P\left(\frac{R_t^{(2n+1)} - nt}{\sqrt{t}} \ge 0\right) - \Phi(0) \gtrsim \frac{1}{\sqrt{t}}, \quad t \ge 1.
$$

We have thus completed the proof of (1.5) for general odd dimensions.

4. Proof of Theorem [1.1](#page-1-2) for even dimensions

4.1. Proof of [\(1.4\)](#page-1-1) for even dimensions. Suppose that $d = 2n$ for **4.1. Proof of (1.4) for even dimensions.** Suppose that $a = 2n$ for some $n \ge 1$. Then for any $t \ge 1$ and $x \ge -(2n-1)\sqrt{t}/2$, we have $T =$ \boldsymbol{x} √ $\bar{t} + (2n - 1)t/2$. We define $J_1^{(2n)}$ $J_1^{(2n)}(t,x)$ and $J_2^{(2n)}$ $2^{(2n)}(t, x)$ as $(3.7)_2$ $(3.7)_2$ and $(3.7)_4$, respectively, with $d = 2n$:

$$
J_1^{(2n)}(t,x)
$$

= $\omega_{2n} \sum_{m=1}^{n-1} \frac{e^{-(2n-1-m)mt/2}}{(2\pi)^m} q_{2n-2m}(t,T) \left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{m-1} \sinh^{2n-2} r \Big|_{r=T}$

and

$$
J_2^{(2n)}(t,x)
$$

= $\omega_{2n} \frac{e^{-n(n-1)t/2}}{(2\pi)^{n-1}} \int_{T}^{\infty} \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{n-1} \sinh^{2n-2} r \right) q_2(t,r) \sinh r \, dr,$

where $J_1^{(2)}(t,x) = 0$ by the convention $\sum_{m=1}^{0} = 0$. In what follows, we simply 1 write $J_1(t, x) = J_1^{(2n)}$ $J_1^{(2n)}(t,x)$ and $J_2(t,x) = J_2^{(2n)}$ $\binom{2n}{2}(t, x)$. We have

(4.1)
$$
P\left(\frac{R_t^{(2n)} - (2n-1)t/2}{\sqrt{t}} \ge x\right) - \Phi(x) = J_1(t, x) + (J_2(t, x) - \Phi(x)).
$$

We now estimate the right hand side of [\(4.1\)](#page-13-1). For an upper bound of $J_1(t, x)$, we can follow the proof of Lemma [3.1](#page-9-6) to show

.

LEMMA 4.1. There exists a constant $c > 0$ such that for any $t \geq 1$,

(4.2)
$$
\sup_{x \ge -(2n-1)\sqrt{t}/2} J_1(t, x) \le \frac{c}{\sqrt{t}}
$$

For $J_2(t, x)$, we have

LEMMA 4.2. There exists a constant $c > 0$ such that for any $t \ge 1$,

(4.3)
$$
\sup_{x \ge -(2n-1)\sqrt{t}/2} |J_2(t,x) - \Phi(x)| \le \frac{c}{\sqrt{t}}.
$$

Proof. Let

$$
a_n(t) = \omega_{2n} \frac{e^{-n(n-1)t/2}}{(2\pi)^{n-1}} = \frac{\pi e^{-n(n-1)t/2}}{2^{n-2}(n-1)!}
$$

and

(4.4)
$$
K(t,x) = \int_{T}^{\infty} \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{n-1} \sinh^{2n-2} r \right) q_2(t,r) \sinh r \, dr.
$$

Then

(4.5)
$$
J_2(t,x) = a_n(t)K(t,x).
$$

By [\(2.2\)](#page-3-5) and the Fubini theorem,

(4.6)
$$
K(t,x) = \frac{2^{1/2}e^{-t/8}}{(2\pi t)^{3/2}} \int_{T}^{x} s e^{-s^2/(2t)} \times \left\{ \int_{T}^{s} \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{n-1} \sinh^{2n-2} r \right) \frac{\sinh r}{(\cosh s - \cosh r)^{1/2}} dr \right\} ds.
$$

By integration by parts,

$$
\int_{T}^{s} \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{n-1} \sinh^{2n-2} r \right) \frac{\sinh r}{(\cosh s - \cosh r)^{1/2}} dr
$$
\n
$$
= \left[-2 \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{n-1} \sinh^{2n-2} r \right) (\cosh s - \cosh r)^{1/2} \right]_{r=T}^{r=s}
$$
\n
$$
+ 2 \int_{T}^{s} \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{n} \sinh^{2n-2} r \right) (\cosh s - \cosh r)^{1/2} \sinh r dr
$$
\n
$$
= 2 \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{n-1} \sinh^{2n-2} r \right) \Big|_{r=T} (\cosh s - \cosh T)^{1/2}
$$
\n
$$
+ 2 \int_{T}^{s} \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{n} \sinh^{2n-2} r \right) (\cosh s - \cosh r)^{1/2} \sinh r dr.
$$

Inductively, we get

$$
(4.7) \int_{T}^{s} \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{n-1} \sinh^{2n-2} r \right) \frac{\sinh r}{(\cosh s - \cosh r)^{1/2}} dr
$$

\n
$$
= \sum_{k=1}^{n-1} \frac{2^{k}}{(2k-1)!!} \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{n-2+k} \right) \sinh^{2n-2} r \Big|_{r=T} (\cosh s - \cosh T)^{k-1/2}
$$

\n
$$
+ \frac{2^{n-1}}{(2n-3)!!} \int_{T}^{s} \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{2n-2} \sinh^{2n-2} r \right) \times (\cosh s - \cosh r)^{n-3/2} \sinh r dr.
$$

Since [\(A.8\)](#page-26-3) yields

$$
\left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{2n-2} \sinh^{2n-2} r = (2n-2)!,
$$

we have

$$
\int_{T}^{s} \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{2n-2} \sinh^{2n-2} r \right) (\cosh s - \cosh r)^{n-3/2} \sinh r \, dr
$$

$$
= (2n-2)! \int_{T}^{s} (\cosh s - \cosh r)^{n-3/2} \sinh r \, dr
$$

$$
= \frac{2(2n-2)!}{2n-1} (\cosh s - \cosh T)^{n-1/2}
$$

and thus

$$
\frac{2^{n-1}}{(2n-3)!!} \int_{T}^{s} \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{2n-2} \sinh^{2n-2} r \right) (\cosh s - \cosh r)^{n-3/2} \sinh r \, dr
$$

$$
= \frac{2^n (2n-2)!}{(2n-1)!!} (\cosh s - \cosh T)^{n-1/2}
$$

$$
= \frac{2^{2n-1} (n-1)!}{2n-1} (\cosh s - \cosh T)^{n-1/2}.
$$

Combining this with [\(4.6\)](#page-14-0) and [\(4.7\)](#page-14-1), we obtain

$$
(4.8) \quad K(t,x) = \frac{2^{1/2}e^{-t/8}}{(2\pi t)^{3/2}} \sum_{k=1}^{n-1} \frac{2^k}{(2k-1)!!} \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{n-2+k} \right) \sinh^{2n-2} r \Big|_{r=T}
$$

$$
\times \int_{T}^{\infty} s e^{-s^2/(2t)} (\cosh s - \cosh T)^{k-1/2} ds
$$

$$
+ \frac{2^{2n-1}(n-1)!}{2n-1} \frac{2^{1/2}e^{-t/8}}{(2\pi t)^{3/2}} \int_{T}^{\infty} s e^{-s^2/(2t)} (\cosh s - \cosh T)^{n-1/2} ds
$$

$$
= K_1(t,x) + K_2(t,x),
$$

whence by (4.5) ,

(4.9)
$$
J_2(t,x) - \Phi(x) = a_n(t)(K_1(t,x) + K_2(t,x)) - \Phi(x)
$$

$$
= a_n(t)K_1(t,x) + (a_n(t)K_2(t,x) - \Phi(x)).
$$

We first estimate $a_n(t)K_1(t, x)$. By [\(A.6\)](#page-26-0), there exists a constant $c_1 > 0$, independent of t, T and k $(1 \le k \le n-1)$, such that

Berry–Esseen bound for the Brownian motions on hyperbolic spaces 17

(4.10)
$$
0 \leq \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{n-2+k} \right) \sinh^{2n-2} r \Big|_{r=1}
$$

$$
\leq c_1 e^{(n-k)T} = c_1 e^{(n-k)(x\sqrt{t} + (2n-1)t/2)}.
$$

On the other hand, for any $s\geq T,$

(4.11)
$$
0 \le (\cosh s - \cosh T)^{k-1/2} \le \cosh^{k-1/2} s \le e^{(k-1/2)s}.
$$

Then by the change of variables $s = u$ $t + (2n - 1)t/2,$

$$
\int_{T}^{\infty} s e^{-s^2/(2t)} (\cosh s - \cosh T)^{k-1/2} ds
$$
\n
$$
\leq \int_{T}^{\infty} s e^{-s^2/(2t)} e^{(k-1/2)s} ds
$$
\n
$$
= \sqrt{t} \int_{x}^{\infty} \left(u \sqrt{t} + \frac{2n-1}{2} t \right) e^{-(u + (2n-1)\sqrt{t}/2)^2/2} e^{(k-1/2)(u\sqrt{t} + (2n-1)t/2)} du
$$
\n
$$
= t \sqrt{t} e^{t/8} e^{-(n^2 - 2nk + k)t/2} \int_{x}^{\infty} \left(\frac{u}{\sqrt{t}} + \frac{2n-1}{2} \right) e^{-u^2/2} e^{-(n-k)u\sqrt{t}} du.
$$

Since

$$
\int_{x}^{\infty} \left(\frac{u}{\sqrt{t}} + \frac{2n-1}{2} \right) e^{-u^2/2} e^{-(n-k)u\sqrt{t}} du
$$

$$
\leq c_2 \int_{x}^{\infty} e^{-(n-k)u\sqrt{t}} du = \frac{c_2 e^{-(n-k)x\sqrt{t}}}{(n-k)\sqrt{t}} \leq \frac{c_2 e^{-(n-k)x\sqrt{t}}}{\sqrt{t}},
$$

we obtain

$$
\int_{T}^{\infty} s e^{-s^2/(2t)} (\cosh s - \cosh T)^{k-1/2} ds
$$

$$
\le c_2 t e^{t/8} e^{-(n^2 - 2nk + k)t/2} e^{-(n-k)x\sqrt{t}},
$$

which yields

$$
\frac{2^{1/2}e^{-t/8}}{(2\pi t)^{3/2}} \int_{T}^{\infty} s e^{-s^2/(2t)} (\cosh s - \cosh T)^{k-1/2} ds
$$

$$
\le c_3 e^{-(n^2-2nk+k)t/2} \frac{e^{-(n-k)x\sqrt{t}}}{\sqrt{t}}.
$$

Combining this with [\(4.10\)](#page-16-0), we get

$$
0 \leq \left(\left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{n-2+k} \right) \sinh^{2n-2} r \Big|_{r=T} \frac{2^{1/2} e^{-t/8}}{(2\pi t)^{3/2}} \times \int_{T}^{\infty} s e^{-s^2/(2t)} (\cosh s - \cosh T)^{k-1/2} ds \leq c_4 e^{(n-k)(x\sqrt{t} + (n-1/2)t)} e^{-(n^2-2nk+k)t/2} \frac{e^{-(n-k)x\sqrt{t}}}{\sqrt{t}} = \frac{c_4 e^{n(n-1)t/2}}{\sqrt{t}}.
$$

This implies that

$$
0 \le K_1(t, x) \le \frac{c_5 e^{n(n-1)t/2}}{\sqrt{t}}
$$

and thus

(4.12)
$$
a_n(t)K_1(t,x) = \frac{\pi e^{-n(n-1)t/2}}{2^{n-2}(n-1)!}K_1(t,x) \le \frac{c_6}{\sqrt{t}}.
$$

Let us next estimate $|a_n(t)K_2(t,x) - \Phi(x)|$. By definition,

(4.13)

$$
a_n(t)K_2(t,x) = b_n(t)\int_T^{\infty} s e^{-s^2/(2t)}(\cosh s - \cosh T)^{n-1/2} ds
$$

= $b_n(t)\int_T^{\infty} \left(s - \frac{2n-1}{2}t\right) e^{-s^2/(2t)}(\cosh s - \cosh T)^{(2n-1)/2} ds$
+ $\frac{2n-1}{2}tb_n(t)\int_T^{\infty} e^{-s^2/(2t)}(\cosh s - \cosh T)^{(2n-1)/2} ds$,

where

(4.14)
$$
b_n(t) = \frac{2^n}{(2n-1)\sqrt{\pi}} \frac{e^{-(2n-1)^2t/8}}{t^{3/2}}.
$$

By the change of variables $y = s/\sqrt{t} - (2n-1)\sqrt{t}/2$ we also have

(4.15)
$$
\Phi(x) = \frac{1}{\sqrt{2\pi}} \int_{x}^{\infty} e^{-y^2/2} dy
$$

$$
= \frac{2n-1}{2} t b_n(t) \int_{T}^{\infty} e^{-s^2/(2t)} \left(\frac{e^s}{2}\right)^{(2n-1)/2} ds.
$$

Then

$$
a_n(t)K_2(t, x) - \Phi(x)
$$

= $b_n(t)\int_T^{\infty} \left(s - \frac{2n-1}{2}t\right)e^{-s^2/(2t)}(\cosh s - \cosh T)^{n-1/2} ds$

$$
- \frac{2n-1}{2}tb_n(t)\int_T^{\infty} e^{-s^2/(2t)} \left\{ \left(\frac{e^s}{2}\right)^{(2n-1)/2} - (\cosh s - \cosh T)^{(2n-1)/2} \right\} ds.
$$

For $s \geq T$, since $e^{s}/2 \geq \cosh s - \cosh T$, we have

$$
(e^s/2)^{(2n-1)/2} \ge (\cosh s - \cosh T)^{(2n-1)/2}.
$$

Therefore,

$$
(4.16) \t |a_n(t)K_2(t,x) - \Phi(x)|
$$

\n
$$
\leq b_n(t) \int_{T}^{\infty} \left| s - \frac{2n-1}{2} t \right| e^{-s^2/(2t)} (\cosh s - \cosh T)^{n-1/2} ds
$$

\n
$$
+ \frac{2n-1}{2} t b_n(t) \int_{T}^{\infty} e^{-s^2/(2t)} \left\{ \left(\frac{e^s}{2} \right)^{(2n-1)/2} - (\cosh s - \cosh T)^{(2n-1)/2} \right\} ds
$$

\n
$$
= b_n(t) L_1(t,x) + \frac{2n-1}{2} t b_n(t) L_2(t,x).
$$

By (4.11) and the change of variables $s = u$ $t + (2n - 1)t/2,$

$$
L_1(t,x) \le \int_{T}^{\infty} \left| s - \frac{2n-1}{2} t \right| e^{-s^2/(2t)} e^{(2n-1)s/2} ds
$$

= $e^{(2n-1)^2 t/8} t \int_{x}^{\infty} |u| e^{-u^2/2} du$
 $\le e^{(2n-1)^2 t/8} t \int_{-\infty}^{\infty} |u| e^{-u^2/2} du = 2e^{(2n-1)^2 t/8} t.$

Hence, by [\(4.14\)](#page-17-0),

(4.17)
$$
b_n(t)L_1(t,x) \leq \frac{2^{n+1}}{(2n-1)\sqrt{\pi}} \frac{1}{\sqrt{t}}.
$$

On the other hand, for any $s\geq T,$

$$
(\cosh s - \cosh T)^{(2n-1)/2} = \cosh^{(2n-1)/2} s \left(1 - \frac{\cosh T}{\cosh s} \right)^{(2n-1)/2}
$$

$$
= \left(\frac{e^s}{2} \right)^{(2n-1)/2} (1 + e^{-2s})^{(2n-1)/2} \left(1 - \frac{\cosh T}{\cosh s} \right)^{(2n-1)/2}
$$

and so

$$
\left(\frac{e^s}{2}\right)^{(2n-1)/2} - (\cosh s - \cosh T)^{(2n-1)/2}
$$
\n
$$
= \left(\frac{e^s}{2}\right)^{(2n-1)/2} \left(1 - (1 + e^{-2s})^{(2n-1)/2} \left(1 - \frac{\cosh T}{\cosh s}\right)^{(2n-1)/2}\right)
$$
\n
$$
\leq \left(\frac{e^s}{2}\right)^{(2n-1)/2} \left(1 - \left(1 - \frac{\cosh T}{\cosh s}\right)^{(2n-1)/2}\right).
$$

Note that for any $p \geq 0$,

(4.18)
$$
0 \le 1 - (1 - v)^p \le pv, \quad 0 \le v \le 1.
$$

Then for any $s \geq T$,

$$
0 \le 1 - \left(1 - \frac{\cosh T}{\cosh s}\right)^{(2n-1)/2} \le \frac{2n-1}{2} \frac{\cosh T}{\cosh s} \le (2n-1)e^{-(s-T)}.
$$

Hence

$$
\left(\frac{e^s}{2}\right)^{(2n-1)/2} - (\cosh s - \cosh T)^{(2n-1)/2} \le (2n-1)e^{(2n-1)s/2}e^{-(s-T)},
$$

which yields

$$
L_2(t,x) = \int_T^{\infty} e^{-s^2/(2t)} \left\{ \left(\frac{e^s}{2}\right)^{(2n-1)/2} - (\cosh s - \cosh T)^{(2n-1)/2} \right\} ds
$$

\n
$$
\leq (2n-1) \int_T^{\infty} e^{-s^2/(2t)} e^{(2n-1)s/2} e^{-(s-T)} ds
$$

\n
$$
= (2n-1)e^{(2n-1)^2t/8} \int_T^{\infty} e^{-(s-(2n-1)t/2)^2/(2t)} e^{-(s-T)} ds
$$

\n
$$
\leq (2n-1)e^{(2n-1)^2t/8} \int_T^{\infty} e^{-(s-T)} ds = (2n-1)e^{(2n-1)^2t/8}.
$$

Therefore, by [\(4.14\)](#page-17-0),

$$
\frac{2n-1}{2}tb_n(t)L_2(t,x) \le \frac{2^{n-1}(2n-1)}{\sqrt{\pi}}\frac{1}{\sqrt{t}}
$$

.

Combining this with [\(4.16\)](#page-18-0) and [\(4.17\)](#page-18-1), for any $t \ge 1$ we obtain

(4.19)
$$
\sup_{x \ge -(2n-1)t/2} |a_n(t)K_2(t,x) - \Phi(x)| \le \frac{2^{n-1}}{\sqrt{\pi}} \left(2n - 1 + \frac{4}{2n-1}\right) \frac{1}{\sqrt{t}}.
$$

This together with [\(4.9\)](#page-15-0) and [\(4.12\)](#page-17-1) leads to [\(4.3\)](#page-13-2). \blacksquare

By (4.2) and (4.3) , we have completed the proof of (1.4) for general even dimensions.

4.2. Proof of [\(1.5\)](#page-1-3) for $d = 2$. We now prove that there exists a constant $c>0$ such that

$$
P\left(\frac{R_t^{(2)} - t/2}{\sqrt{t}} \ge 0\right) - \Phi(0) \ge \frac{c}{\sqrt{t}}, \quad t \ge 1.
$$

Let $R_t = R_t^{(2)}$ $t^{(2)}$. Then by [\(2.1\)](#page-3-4) with $\omega_2 = 2$,

$$
P\left(\frac{R_t - t/2}{\sqrt{t}} \ge 0\right) = P\left(R_t \ge \frac{t}{2}\right) = 2\int_{t/2}^{\infty} q_2(t, r) \sinh r \, dr.
$$

On the other hand, by [\(4.4\)](#page-14-3) and [\(4.8\)](#page-15-1) with $n = 1$ and $x = 0$,

$$
2\int_{t/2}^{\infty} q_2(t,r) \sinh r \, dr = \frac{2e^{-t/8}}{t\sqrt{\pi t}} \int_{t/2}^{\infty} r e^{-r^2/(2t)} (\cosh r - \cosh(t/2))^{1/2} \, dr.
$$

Moreover, by integration by parts,

$$
\frac{2e^{-t/8}}{t\sqrt{\pi t}}\int_{t/2}^{\infty}re^{-r^2/(2t)}(\cosh r - \cosh(t/2))^{1/2} dr
$$

=
$$
\frac{e^{-t/8}}{\sqrt{\pi t}}\int_{t/2}^{\infty}e^{-r^2/(2t)}\frac{\sinh r}{(\cosh r - \cosh(t/2))^{1/2}} dr.
$$

Therefore,

$$
P\left(\frac{R_t - t/2}{\sqrt{t}} \ge 0\right) = \frac{e^{-t/8}}{\sqrt{2\pi t}} \int_{t/2}^{\infty} e^{-r^2/(2t)} \frac{\sqrt{2}\sinh r}{(\cosh r - \cosh(t/2))^{1/2}} dr.
$$

On the other hand, by [\(4.15\)](#page-17-2) with $x = 0$ and $n = 1$,

$$
\Phi(0) = \frac{e^{-t/8}}{\sqrt{2\pi t}} \int_{t/2}^{\infty} e^{-r^2/(2t)} e^{r/2} dr.
$$

We thus have

$$
(4.20) \qquad P\left(\frac{R_t - t/2}{\sqrt{t}} \ge 0\right) - \Phi(0)
$$
\n
$$
= \frac{e^{-t/8}}{\sqrt{2\pi t}} \int_{t/2}^{\infty} e^{-r^2/(2t)} e^{r/2} \left(\frac{\sqrt{2}e^{-r/2}\sinh r}{(\cosh r - \cosh(t/2))^{1/2}} - 1\right) dr
$$
\n
$$
= \frac{e^{-t/8}}{\sqrt{2\pi t}} \int_{t/2}^{\infty} e^{-r^2/(2t)} e^{r/2} F_t(r) dr
$$

with

$$
F_t(r) = \frac{\sqrt{2} e^{-r/2} \sinh r}{(\cosh r - \cosh(t/2))^{1/2}} - 1.
$$

Then $F_t(r) \geq 0$ for any $t \geq 1$ and $r \geq t/2$ because

$$
(4.21) \quad (\sqrt{2}e^{-r/2}\sinh r)^2 - (\cosh r - \cosh(t/2))
$$

= $\frac{e^{-3r} - 3e^{-r}}{2} + \frac{e^{t/2} + e^{-t/2}}{2}$
 $\ge -\frac{3e^{-t/2}}{2} + \frac{e^{t/2} + e^{-t/2}}{2} = \frac{e^{t/2} - 2e^{-t/2}}{2} \ge 0.$

Hence the right hand side of [\(4.20\)](#page-20-0) is positive.

For $t\geq 1$ and $r\geq t/2,$ let

(4.22)
$$
F_t(r) = \frac{\sqrt{2} e^{-r/2} \sinh r}{\cosh^{1/2} r} \left(1 - \frac{\cosh(t/2)}{\cosh r} \right)^{-1/2} - 1
$$

$$
= \left(1 - \frac{\cosh(t/2)}{\cosh r} \right)^{-1/2} - 1 - \left(1 - \frac{\cosh(t/2)}{\cosh r} \right)^{-1/2} \left(1 - \frac{\sqrt{2} e^{-r/2} \sinh r}{\cosh^{1/2} r} \right).
$$

Since

$$
\frac{1}{\sqrt{1-v}} - 1 \ge 1 - \sqrt{1-v} \ge \frac{1}{2}v, \quad 0 \le v < 1,
$$

we obtain

(4.23)
$$
\left(1 - \frac{\cosh(t/2)}{\cosh r}\right)^{-1/2} - 1 \ge \frac{1}{2} \frac{\cosh(t/2)}{\cosh r} \ge \frac{1}{4} e^{-(r-t/2)}.
$$

We also note that

$$
0 \le 1 - \frac{1 - v}{\sqrt{1 + v}} = \frac{3v - v^2}{\sqrt{1 + v}(\sqrt{1 + v} + 1 - v)} \le 3v, \quad 0 \le v \le 1,
$$

and so

$$
(4.24)
$$

$$
0 \le 1 - \frac{\sqrt{2}e^{-r/2}\sinh r}{\cosh^{1/2} r} = 1 - \frac{e^{r/2} - e^{-3r/2}}{\sqrt{e^r + e^{-r}}} = 1 - \frac{1 - e^{-2r}}{\sqrt{1 + e^{-2r}}} \le 3e^{-2r}.
$$

If we assume in addition that $r \geq 1 + t/2$, then

$$
1 - \frac{\cosh(t/2)}{\cosh r} \ge 1 - \frac{\cosh(t/2)}{\cosh(1 + t/2)} \ge 1 - \frac{1}{e} > 0.
$$

Combining this with [\(4.24\)](#page-21-0) we have, for any $r \ge 1 + t/2$,

$$
(4.25)\quad \left(1 - \frac{\cosh(t/2)}{\cosh r}\right)^{-1/2} \left(1 - \frac{\sqrt{2}e^{-r/2}\sinh r}{\cosh^{1/2}r}\right) \le 3(1 - e^{-1})^{-1/2}e^{-2r}.
$$

Therefore, by [\(4.22\)](#page-21-1), [\(4.23\)](#page-21-2) and [\(4.25\)](#page-21-3), for any $r \ge 1 + t/2$ we have

(4.26)
$$
F_t(r) \geq \frac{1}{4}e^{-(r-t/2)} - 3(1 - e^{-1})^{-1/2}e^{-2r}.
$$

By (4.26) , we get

$$
(4.27) \int_{t/2}^{\infty} e^{-r^2/(2t)} e^{r/2} F_t(r) \, \mathrm{d}u \ge \int_{1+t/2}^{\infty} e^{-r^2/(2t)} e^{r/2} F_t(r) \, \mathrm{d}r
$$

$$
\ge \frac{1}{4} e^{t/2} \int_{1+t/2}^{\infty} e^{-r^2/(2t) - r/2} \, \mathrm{d}r
$$

$$
- 3(1 - e^{-1})^{-1/2} \int_{1+t/2}^{\infty} e^{-r^2/(2t) - 3r/2} \, \mathrm{d}r.
$$

Note that for any $c > 0$, by (2.8) we have

$$
(4.28) \int_{1+t/2}^{\infty} e^{-r^2/(2t)} e^{-cr} dr
$$

= $e^{c^2t/2} \int_{1+t/2}^{\infty} e^{-(r+ct)^2/(2t)} dr = e^{c^2t/2} \sqrt{t} \int_{(c+1/2)\sqrt{t}+1/\sqrt{t}}^{\infty} e^{-u^2/2} du$
 $\asymp e^{c^2t/2} e^{-(c+1/2)^2t/2} = e^{-(4c+1)t/8}.$

At the second equality above, we have used the change of variables $u =$ $(r + ct)/\sqrt{t}$. Hence by [\(4.27\)](#page-22-1), and taking $c = 1/2$ and $c = 3/2$ in [\(4.28\)](#page-22-2), for all large $t \geq 1$ we get

(4.29)
$$
\int_{t/2}^{\infty} e^{-r^2/(2t)} e^{r/2} F_t(r) \, \mathrm{d}u \ge c_1 e^{t/2} \cdot e^{-3t/8} - c_2 e^{-7t/8} \ge c_3 e^{t/8}.
$$

Recall that the integral above is positive for any $t \ge 1$. By [\(4.20\)](#page-20-0) and [\(4.29\)](#page-22-3) we then have

$$
P\left(\frac{R_t - t/2}{\sqrt{t}} \ge 0\right) - \Phi(0) \gtrsim \frac{1}{\sqrt{t}}, \quad t \ge 1.
$$

Thus, we have arrived at the desired conclusion.

REMARK 4.3. At the moment, it is unknown whether (1.5) is valid for even $d \geq 4$. To explain this, let us apply the argument for $d = 2$ to general even dimension $d = 2n$ with $n \ge 2$. Fix $x \ge 0$. Then by [\(4.1\)](#page-13-1) and [\(4.9\)](#page-15-0),

(4.30)
\n
$$
P\left(\frac{R_t^{(2n)} - (2n - 1)t/2}{\sqrt{t}} \ge x\right) - \Phi(x) = J_1(t, x) + (J_2(t, x) - \Phi(x))
$$
\n
$$
= J_1(t, x) + a_n(t)K_1(t, x) + (a_n(t)K_2(t, x) - \Phi(x)).
$$

By the expression in the first line of [\(4.13\)](#page-17-3), and integration by parts,

$$
a_n(t)K_2(t,x) = b_n(t)\int_T^{\infty} s e^{-s^2/(2t)}(\cosh s - \cosh T)^{n-1/2} ds
$$

=
$$
\frac{2n-1}{2}tb_n(t)\int_T^{\infty} e^{-s^2/(2t)} \sinh s (\cosh s - \cosh T)^{n-3/2} ds.
$$

Hence if we let

(4.31)
$$
G_t(s) = 1 - \left(\frac{2}{e^s}\right)^{(2n-1)/2} \sinh s \left(\cosh s - \cosh T\right)^{n-3/2}
$$

$$
= 1 - (1 - e^{-2s})(1 + e^{-2s} - 2e^{-s}\cosh T)^{n-3/2}, \quad s \ge T,
$$

then by (4.15) ,

$$
a_n(t)K_2(t,x) - \Phi(x) = -\frac{2n-1}{2}tb_n(t)\int_T^{\infty} e^{-s^2/(2t)} \left(\frac{e^s}{2}\right)^{(2n-1)/2} G_t(s) ds.
$$

If $n = 1$, then $G_t(s)$ is negative by a similar calculation to [\(4.21\)](#page-21-5) and so the right hand side above is positive. On the other hand, if $n \geq 2$, then we can see by [\(4.31\)](#page-23-0) that the function $G_t(s)$ is positive and decreasing on $[T, \infty)$, and such that $G_t(T) = 1$ and $G_t(s) \to 0$ as $s \to \infty$. Therefore, [\(4.19\)](#page-19-0) implies that for any $t \geq 1$,

$$
a_n(t)K_2(t,x) - \Phi(x) \ge -|a_n(t)K_2(t,x) - \Phi(x)|
$$

$$
\ge -\frac{2^{n-1}}{\sqrt{\pi}} \left(2n - 1 + \frac{4}{2n - 1}\right) \frac{1}{\sqrt{t}} = -\frac{c_*}{\sqrt{t}}.
$$

We also see that, by an argument similar to that for (3.19) and (4.12) together with $(A.7)$, there exist positive constants c_1 and c_2 such that for any $t \geq 1$,

$$
J_1(t,x) \ge \frac{c_1}{\sqrt{t}}, \quad a_n(t)K_1(t,x) \ge \frac{c_2}{\sqrt{t}}.
$$

Hence [\(4.30\)](#page-22-4) yields, for any $t \geq 1$,

$$
P\left(\frac{R_t^{(2n)} - (2n - 1)t/2}{\sqrt{t}} \ge x\right) - \Phi(x) \ge \frac{c_1 + c_2 - c_*}{\sqrt{t}}.
$$

However, since we could not get effective estimates for c_1 and c_2 , it is unknown whether the right hand side above is positive or not. We note that the constants c_1 and c_2 come from [\(2.5\)](#page-3-1) and [\(A.7\)](#page-26-2), respectively.

A. Appendix. To prove Theorem [1.1,](#page-1-2) we establish the following two lemmas.

LEMMA A.1. Let $n \geq 1$ and $r \neq 0$.

(i) For any $k = 1, 2, \ldots$,

(A.1)
$$
\left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{2k} \sinh^{2n-1} r\n= \sum_{l=0}^{k} \frac{(2k)!}{2^{k-l}(k-l)!(2l)!} \left(\prod_{m=0}^{k+l-1} (2n-1-2m)\right)\n\times \cosh^{2l} r \sinh^{2n-1-(2k+2l)} r.
$$

For any $k = 0, 1, \ldots$,

(A.2)
$$
\left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{2k+1} \sinh^{2n-1} r = \sum_{l=1}^{k+1} \frac{(2k+1)!}{2^{k+1-l}(k+1-l)!(2l-1)!}
$$

$$
\times \left(\prod_{m=0}^{k+l-1} (2n-1-2m)\right) \cosh^{2l-1} r \sinh^{2n-1-(2k+2l)} r.
$$

(ii) For any $k = 1, \ldots, n$,

(A.3)
$$
\left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{2k} \sinh^{2n} r
$$

=
$$
\sum_{l=0}^{k \wedge (n-k)} \frac{(2k)!}{2^{k-l}(k-l)!(2l)!} \left(\prod_{m=0}^{k+l-1} (2n-2m)\right) \cosh^{2l} r \sinh^{2n-(2k+2l)} r.
$$

For any $k = 0, 1, ..., n - 1$,

(A.4)
$$
\left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{2k+1} \sinh^{2n} r = \sum_{l=1}^{(k+1)\wedge (n-k)} \frac{(2k+1)!}{2^{k+1-l}(k+1-l)!(2l-1)!}
$$

$$
\times \left(\prod_{m=0}^{k+l-1} (2n-2m)\right) \cosh^{2l-1} r \sinh^{2n-(2k+2l)} r.
$$

Proof. Fix $n \geq 1$. Since the proof of (i) is similar to that of (ii), we present the former only.

We prove $(A.1)$ by induction. For $k = 1$, we have

$$
\left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^2 \sinh^{2n-1} r
$$

= $(2n - 1) \sinh^{2n-3} r + (2n - 1)(2n - 3) \cosh^2 r \sinh^{2n-5} r$,

whence [\(A.1\)](#page-24-0) is valid for $k = 1$. Suppose that it holds for some $k \ge 1$. Since

$$
\frac{\partial}{\partial r} (\cosh^{2l} r \sinh^{2n-1-(2k+2l)} r)
$$

= $2l \cosh^{2l-1} r \sinh^{2n-(2k+2l)} r$
+ $(2n - 1 - (2k + 2l)) \cosh^{2l+1} r \sinh^{2n-2-(2k+2l)} r$,

we obtain

$$
(A.5) \frac{\partial}{\partial r} \left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{2k} \sinh^{2n-1} r
$$
\n
$$
= \sum_{l=0}^{k} \frac{(2k)!}{2^{k-l}(k-l)!(2l)!} \left(\prod_{m=0}^{k+l-1} (2n-1-2m) \right) \frac{\partial}{\partial r} (\cosh^{2l} r \sinh^{2n-1-(2k+2l)} r)
$$
\n
$$
= \sum_{l=1}^{k} \frac{(2k)!}{2^{k-l}(k-l)!(2l-1)!} \left(\prod_{m=0}^{k+l-1} (2n-1-2m) \right) \cosh^{2l-1} r \sinh^{2n-(2k+2l)} r
$$
\n
$$
+ \sum_{l=0}^{k} \frac{(2k)!}{2^{k-l}(k-l)!(2l)!} \left(\prod_{m=0}^{k+l} (2n-1-2m) \right) \cosh^{2l+1} r \sinh^{2n-2-(2k+2l)} r
$$
\n
$$
= \sum_{l=1}^{k} \frac{(2k)!}{2^{k-l}(k-l)!(2l-1)!} \left(\prod_{m=0}^{k+l-1} (2n-1-2m) \right) \cosh^{2l-1} r \sinh^{2n-(2k+2l)} r
$$
\n
$$
+ \sum_{l=1}^{k+1} \frac{(2k)!}{2^{k-l+1}(k-l+1)!(2l-2)!}
$$
\n
$$
\times \left(\prod_{m=0}^{k+l-1} (2n-1-2m) \right) \cosh^{2l-1} r \sinh^{2n-(2k+2l)} r
$$
\n
$$
= \sum_{l=1}^{k+1} \frac{(2k+1)!}{2^{k+l-l}(k+1-l)!(2l-1)!} \left(\prod_{m=0}^{k+l-1} (2n-1-2m) \right)
$$
\n
$$
\times \cosh^{2l-1} r \sinh^{2n-(2k+2l)} r.
$$

Therefore, [\(A.2\)](#page-24-1) holds. Using it we obtain, in a similar way to [\(A.5\)](#page-25-0),

$$
\frac{\partial}{\partial r} \left(\frac{1}{\sinh r} \frac{\partial}{\partial r} \right)^{2k+1} \sinh^{2n-1} r\n= \sum_{l=0}^{k+1} \frac{(2k+2)!}{2^{k+1-l}(k+1-l)!(2l)!} \left(\prod_{m=0}^{k+l} (2n-1-2m) \right) \cosh^{2l} r \sinh^{2n-2-(2k+2l)} r
$$

and thus

$$
\left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{2k+2} \sinh^{2n-1} r
$$

=
$$
\sum_{l=0}^{k+1} \frac{(2k+2)!}{2^{k+1-l}(k+1-l)!(2l)!} \left(\prod_{m=0}^{k+l} (2n-1-2m)\right) \cosh^{2l} r \sinh^{2n-3-(2k+2l)} r.
$$

That is, $(A.1)$ is valid for k replaced with $k + 1$. Hence the induction is complete. In particular, our argument above also implies $(A.2)$ for $k \geq 1$. Since we can verify $(A.2)$ for $k = 0$ by direct computation, the proof is complete. \blacksquare

REMARK A.2. Let m and n be positive integers with $n \geq 2m$. Then by Lemma [A.1,](#page-24-2) we have the following two clauses:

• There exists a constant $c_1 > 0$ such that

(A.6)
$$
\left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^m \sinh^n r \leq c_1 e^{(n-m)r}, \quad r > 0.
$$

Note that for even *n* we see from $(A.3)$ and $(A.4)$ that $(A.6)$ is still valid for $n \geq m$.

• For any $r_0 > 0$, there exists a constant $c_2 > 0$ such that

(A.7)
$$
\left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^m \sinh^n r \ge c_2 e^{(n-m)r}, \quad r \ge r_0.
$$

LEMMA A.3. Let $r \neq 0$.

(i) For $n \geq 1$,

(A.8)
$$
\left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{2n} \sinh^{2n} r = (2n)!
$$

(ii) For $n \geq 1$,

(A.9)
$$
\left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^n \sinh^{2n+1} r = \frac{(2n+1)!!}{n+1} \sinh((n+1)r).
$$

Proof. (i) follows from $(A.3)$. To show (ii), we calculate the left hand side of $(A.9)$ for odd n and even n separately. Namely, we will prove that for any $k \geq 1$,

(A.10)
$$
\left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{2k-1} \sinh^{4k-1} r = \frac{(4k-1)!!}{2k} \sinh(2kr)
$$

and

(A.11)
$$
\left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{2k} \sinh^{4k+1} r = \frac{(4k+1)!!}{2k+1} \sinh((2k+1)r).
$$

By (A.2),
\n
$$
\left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{2k-1} \sinh^{4k-1} r
$$
\n
$$
= \sum_{l=1}^{k} \frac{(2k-1)!}{2^{k-l}(k-l)!(2l-1)!} \left(\prod_{m=0}^{k+l-2} (4k-1-2m)\right) \cosh^{2l-1} r \sinh^{2k-2l+1} r
$$
\n
$$
= \sum_{j=0}^{k-1} \frac{(2k-1)!}{2^j j!(2k-2j-1)!} \left(\prod_{m=0}^{2k-j-2} (4k-1-2m)\right) \cosh^{2k-(2j+1)} r \sinh^{2j+1} r,
$$

where we set $j = k - l$ in the second equality above. If $0 \le j \le k - 1$, then

$$
\frac{(2k-1)!}{2^j j! (2k-2j-1)!} \prod_{m=0}^{2k-j-2} (4k-1-2m) = \frac{(2k-1)!}{(2j)!! (2k-2j-1)!} \prod_{l=j+2}^{2k} (2l-1)
$$

$$
= \frac{(4k-1)!!}{2k} \frac{(2k)!}{(2j+1)!(2k-(2j+1))!} = \frac{(4k-1)!!}{2k} {2k \choose 2j+1}
$$

and thus

$$
\left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{2k-1} \sinh^{4k-1} r
$$

=
$$
\frac{(4k-1)!!}{2k} \sum_{j=0}^{k-1} {2k \choose 2j+1} \cosh^{2k-(2j+1)} r \sinh^{2j+1} r = \frac{(4k-1)!!}{2k} \sinh 2kr.
$$

The proof of [\(A.10\)](#page-26-4) is complete.

By [\(A.1\)](#page-24-0) and a similar argument to the above,

$$
\left(\frac{1}{\sinh r} \frac{\partial}{\partial r}\right)^{2k} \sinh^{4k+1} r
$$
\n
$$
= \sum_{l=0}^{k} \frac{(2k)!}{2^{k-l}(k-l)!(2l)!} \left(\prod_{m=0}^{k+l-1} (4k+1-2m)\right) \cosh^{2l} r \sinh^{2k-(2l-1)} r
$$
\n
$$
= \sum_{j=0}^{k} \frac{(2k)!}{2^j j!(2(k-j))!} \left(\prod_{i=j+1}^{2k} (2i+1)\right) \cosh^{2(k-j)} r \sinh^{2j+1} r
$$
\n
$$
= \frac{(4k+1)!!}{2k+1} \sum_{j=0}^{k} {2k+1 \choose 2j+1} \cosh^{2(k-j)} r \sinh^{2j+1} r
$$
\n
$$
= \frac{(4k+1)!!}{2k+1} \sinh((2k+1)r).
$$

The proof of $(A.11)$ is complete.

Acknowledgements. The author is grateful to the referee for his/her careful reading of the manuscript, and for helpful comments and suggestions, which simplified the argument and improved the presentation. This work was supported by JSPS KAKENHI Grant Numbers JP22K18675, JP23H01076, JP23K25773.

References

- [\[1\]](http://dx.doi.org/10.1016/0022-1236(92)90133-4) J.-Ph. Anker and A. G. Setti, Asymptotic finite propagation speed for heat diffusion on certain Riemannian manifolds, J. Funct. Anal. 103 (1992), 50–61.
- [\[2\]](http://dx.doi.org/10.1007/BF02213570) M. Babillot, A probabilistic approach to heat diffusion on symmetric spaces, J. Theoret. Probab. 7 (1994), 599–607.
- [\[3\]](http://dx.doi.org/10.1007/s10955-014-0939-5) V. Cammarota, A. De Gregorio and C. Macci, On the asymptotic behavior of the hyperbolic Brownian motion, J. Statist. Phys. 154 (2014), 1550–1568.
- [\[4\]](http://dx.doi.org/10.1017/CBO9780511566158) E. B. Davies, *Heat Kernels and Spectral Theory*, Cambridge Univ. Press, Cambridge, 1989.
- [\[5\]](http://dx.doi.org/10.1112/plms/s3-57.1.182) E. B. Davies and N. Mandouvalos, Heat kernel bounds on hyperbolic space and Kleinian groups, Proc. London Math. Soc. 57 (1988), 182–208.
- [6] R. Durrett, Probability: Theory and Examples, 5th ed., Cambridge Univ. Press, Cambridge, 2019.
- [\[7\]](http://dx.doi.org/10.1007/978-0-387-85650-6_10) A. Grigor'yan and E. Hsu, Volume growth and escape rate of Brownian motion on a Cartan–Hadamard manifold, in: Sobolev Spaces in Mathematics II, Int. Math. Ser. (N.Y.) 9, Springer, New York, 2009, 209–225.
- [\[8\]](http://dx.doi.org/10.1112/S0024609398004780) A. Grigor'yan and M. Noguchi, The heat kernel on hyperbolic space, Bull. London Math. Soc. 30 (1998), 643–650.
- [\[9\]](http://dx.doi.org/10.1080/17442509608834035) J.-C. Gruet, Semi-groupe du mouvement brownien hyperbolique, Stochast. Stochast. Rep. 56 (1996), 53–61.
- [\[10\]](http://dx.doi.org/10.1016/j.spl.2011.06.005) M. Hirao, Large deviations for the radial processes of the Brownian motions on hyperbolic spaces, Statist. Probab. Lett. 81 (2011), 1561–1564.
- [\[11\]](http://dx.doi.org/10.1090/gsm/038) E. P. Hsu, Stochastic Analysis on Manifolds, Amer. Math. Soc., Providence, RI, 2002.
- [12] I. Karatzas and S. E. Shreve, Brownian Motion and Stochastic Calculus, 2nd ed., Springer, New York, 1991.
- [\[13\]](http://dx.doi.org/10.4064/cm119-2-3) H. Matsumoto, Limiting behaviors of the Brownian motions on hyperbolic spaces, Colloq. Math. 119 (2010), 193–215.
- [\[14\]](http://dx.doi.org/10.3792/pjaa.93.27) Y. Shiozawa, Escape rate of the Brownian motions on hyperbolic spaces, Proc. Japan Acad. Ser. A Math. Sci. 93 (2017), 27–29.
- [\[15\]](http://dx.doi.org/10.4310/CAG.2022.v30.n9.a7) J. L. Vázquez, Asymptotic behaviour for the heat equation in hyperbolic space, Comm. Anal. Geom. 30 (2022), 2123–2156.

Yuichi Shiozawa Department of Mathematical Sciences Faculty of Science and Engineering Doshisha University Kyotanabe, Kyoto, 610-0394, Japan E-mail: yshiozaw@mail.doshisha.ac.jp